# Electroweak Precision at CHF and SppC

Testing the consistence of the SM
 Indirect search of NP beyond the SM

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### Top discovery: EW theory tests at Loop level



### W-boson, Top-quark and Higgs boson

• Highly correlated at the quantum level



## Top quark and 125GeV Higgs boson





### Higgs boson couplings

• New set of reference SM parameters  $m_H \sim 126 \text{ GeV}$   $\Gamma_H = 4.2 \text{ MeV}$   $\lambda = (m_H/v)^2/2 = 0.131$ 

$Br(H \to WW^*) = 23\%$	*	-2lnλ(μ)<1 Intervals	2011 - 2012 Data
$Br(H \to ZZ^*) = 2.9\%$	*	ATLAS Preliminary W,Z H → bb √s = 7 TeV: ∫Ldt = 4.6-4.7 fb <sup>-1</sup>	
$Br(H \to bb) = 56\%$	*	$H \rightarrow \tau\tau$ $\sqrt{s} = 7 \text{ TeV}: \int Ldt = 4.7 \text{ fb}^{-1}$ $H \rightarrow WW^{(*)} \rightarrow IvIv$	
$Br(H \to cc) = 2.8\%$		$\sqrt{s} = 7 \text{ TeV}: \int Ldt = 4.7 \text{ fb}^{-1}$ $H \longrightarrow \gamma \gamma$ $\sqrt{s} = 8 \text{ TeV}: (1 dt = 5.9 \text{ fb}^{-1})$	
$Br(H \to \tau \tau) = 6.2\%$	*	$\sqrt{s} = 7 \text{ TeV}: \int Ldt = 4.8 \text{ fb}^{-1}$ H $\rightarrow ZZ^{\binom{n}{2}} \rightarrow IIII$	
$Br(H \to \mu\mu) = 0.021\%$		√s = 8 TeV: ∫Ldt = 5.8 fb <sup>-1</sup> √s = 7 TeV: ∫Ldt = 4.8 fb <sup>-1</sup>	-
$Br(H \to gg) = 8.5\%$		Combined (8 = 8 TeV: fLdt = 5.8 - 5.9 fb <sup>-1</sup> (8 = 7 TeV: fLdt = 4.6 - 4.8 fb <sup>-1</sup>	• $\mu = 1.2^{+0.3}_{-0.3}$
$Br(H \to \gamma \gamma) = 0.23\%$	*	-1	0 1
$\operatorname{Br}(H \to \gamma Z) = 0.16\%$	*	S	ignal strength (μ)

### Higgs boson couplings at LC Peskin, 1208.5152







### e<sup>+</sup>e<sup>-</sup> collider at 250 GeV

- The LHC 7-8TeV results imply no need for a LEP above 500 GeV.
- If the simple scalar Higgs model is correct, the Higgs couplings to each particle is proportional to its mass.

We can test this hypothesis to high accuracy.



### Higgs effective coupling



### CHF at 250 GeV

Test the notorious
 3 sigma deviation
 in the A<sub>FB</sub> of
 bottom quark

 Measure threegauge-boson coupling

7 effective couplings

	Measurement	Pull	Pull -3 -2 -1 0 1 2 3
m <sub>z</sub> [GeV]	91.1875 ± 0.0021	.04	
Г <sub>Z</sub> [GeV]	2.4952 ± 0.0023	46	-
$\sigma_{hadr}^{0}\left[nb ight]$	41.540 ± 0.037	1.62	
R <sub>I</sub>	20.767 ± 0.025	1.09	
A <sup>0,I</sup> <sub>fb</sub>	0.01714 ± 0.00095	.79	-
A <sub>e</sub>	0.1498 ± 0.0048	.41	-
Α <sub>τ</sub>	0.1439 ± 0.0041	96	-
$sin^2 \theta_{eff}^{lept}$	0.2322 ± 0.0010	.78	-
m <sub>w</sub> [GeV]	80.446 ± 0.040	1.32	
R <sub>b</sub>	0.21664 ± 0.00068	1.32	
R <sub>c</sub>	0.1729 ± 0.0032	.20	•
A <sup>0,b</sup>	0.0982 ± 0.0017	-3.20	
A <sup>0,c</sup> <sub>fb</sub>	0.0689 ± 0.0035	-1.48	
A <sub>b</sub>	0.921 ± 0.020	68	-
A <sub>c</sub>	$0.667 \pm 0.026$	05	
A <sub>l</sub>	0.1513 ± 0.0021	1.68	
sin <sup>2</sup> θ <sub>W</sub>	0.2255 ± 0.0021	1.20	
m <sub>w</sub> [GeV]	80.452 ± 0.062	.95	
m <sub>t</sub> [GeV]	174.3 ± 5.1	27	-
$\Delta lpha_{had}^{(5)}(m_Z)$	0.02761 ± 0.00036	36	
			-3 -2 -1 0 1 2 3

# 50 TeV versus 14 TeV

### Parton distribution function



Gluon induced channels are highly suppressed, while the quark-antiquark channels are less suppressed.

 For heavy resonance production, the quark-(anti)quark initial states dominate.

### Proton-Proton at 50<sup>+</sup> TeV



The cross sections of quark-quark initial state increase by a factor of 3-5 while the cross sections of gluon-gluon initial state increase by a factor 5-10.

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### NLO QCD corrections to heavy quark production

• QQ production via the QCD Interaction



#### PDF uncertainties

Berger and QHC, Phys. Rev. D81 (2010) 035006

## Pros and Cons of 50 TeV SppC

 The effective (x) is lowered by a factor of 3.5 when machine energy increases from 14TeV to 50TeV. For a TeV resonance,

 $\langle x \rangle_{14} \sim \frac{\text{TeV}}{14 \text{ TeV}} \sim 0.07$   $\langle x \rangle_{50} \sim \frac{\text{TeV}}{50 \text{ TeV}} \sim 0.02$ The gluon PDF exhibits a larger uncertainty.

New Data  $\longrightarrow$  Proton structure in small x

The cross section of New physics resonance (in the large \large \large x \rangle region) production increases less than the cross section of the SM backgrounds (in the small \large x \rangle region).



# What is not measured yet • One last not-measured fermion gauge coupling $g_Z t \bar{t}$



#### It always comes together with W-t-b coupling.

• Neutrino: Dirac or Majorana

• Higgs self interaction coupling

### EWPT: Bottom-Up approach

- Effective Field Theory
  - Gauge invariant (less model independent)
  - Easy to track the origin and order of NP
  - Too many operators

- Effective Lagrangian (effective coupling)
  More general (Lorentz invariant)
  - Tree and loop effects messed up.

### Tree-level induced dim-6 operators

$$\begin{array}{rcl}
\mathcal{O}_{\phi q}^{(1)} &=& i \left(\phi^{\dagger} D_{\mu} \phi\right) \left(\bar{q} \gamma^{\mu} q\right), \\
\mathcal{O}_{\phi q}^{(3)} &=& i \left(\phi^{\dagger} \tau^{I} D_{\mu} \phi\right) \left(\bar{q} \gamma^{\mu} \tau^{I} q\right), \\
\mathcal{O}_{\phi t} &=& i \left(\phi^{\dagger} D_{\mu} \phi\right) \left(\bar{t}_{R} \gamma^{\mu} t_{R}\right), \\
\mathcal{O}_{\phi b} &=& i \left(\phi^{\dagger} D_{\mu} \phi\right) \left(\bar{b}_{R} \gamma^{\mu} b_{R}\right), \\
\mathcal{O}_{\phi \phi} &=& \left(\phi^{\dagger} \epsilon D_{\mu} \phi\right) \left(\bar{t}_{R} \gamma^{\mu} b_{R}\right), \\
\end{array}$$



### Effective wtb, ztt and zbb couplings



### Effective Wtb, Ztt, Zbb couplings

• New parameterization of couplings

$$\mathcal{O}_{Wtb} = \frac{g}{\sqrt{2}} \mathcal{F}_L W^+_\mu \bar{t}_L \gamma^\mu b_L + h.c. ,$$
  
$$\mathcal{O}_{Zt\bar{t}} = \frac{g}{2c_w} Z_\mu \left( 2\mathcal{F}_L \bar{t}_L \gamma^\mu t_L + \mathcal{F}_R \bar{t}_R \gamma^\mu t_R \right)$$

• The coefficients of the left-handed neutral and charged currents are related,

which is predicted by the EW gauge symmetry after the stringent constraint on  $Zb_Lb_L$  imposed.

### How to probe such a correlation

$$\mathcal{O}_{Wtb} = \frac{g}{\sqrt{2}} \mathcal{F}_L W^+_\mu \bar{t}_L \gamma^\mu b_L + h.c. ,$$
  
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• At the Hadron Collider

 U. Baur, A. Juste, L.H. Orr, D. Rainwater Phys.Rev.D71:054013,2005; Phys. Rev.D73:034016,2006

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#### • At the Linear Collider

P. Batra, T. Tait, Phys.Rev.D74:054021,2006 QHC, J. Wudka, Phys.Rev.D74: 094015, 2006

 $\overline{b}$ 



### Impact of anomalous couplings on $\sigma_t$ and $\sigma_{t\bar{t}}$

• Inclusive cross sections of single-t and Ztt productions:

$$\sigma_t = \sigma_t^0 \left[ 1 + 2\mathcal{F}_L + 2\delta V_{tb} + \mathcal{O}\left(\mathcal{F}_L^2, \delta V_{tb}^2\right) \right],$$
  
$$\sigma_{Zt\bar{t}} = \sigma_{Zt\bar{t}}^0 \left[ 1 + 4.4\mathcal{F}_L - 1.5\mathcal{F}_R + \mathcal{O}\left(\mathcal{F}_L^2, \mathcal{F}_R^2, \mathcal{F}_L\mathcal{F}_R\right) \right]$$
  
$$\delta\sigma = (\sigma - \sigma^0)/\sigma^0 \quad \delta V_{tb} = |V_{tb}|^{(\text{exp})} - |V_{tb}|^{(\text{SM})}$$

 $\delta V_{tb} = -0.23\delta\sigma_{Zt\bar{t}} + 0.5\delta\sigma_t - 0.34\mathcal{F}_R$ 

Ed L. Berger, QHC, Ian Low, Phys.Rev.D80: 074020 (2009)

Note: Vtb cannot be extracted out from single top production alone.

### Experiments versus Theories

• Physics is associated with many scales

