Research on Cosmic microwave background Radiation

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Institute of High Energy Physics Since 1980s



Dayabay neutrino experiment

Beijing Spectrometer III

Particle Astrophysics Division

- Hard X-ray satellite
- Cosmic ray experiment
- > AliCPT in Tibet



Outline

- What is CMB ?
- Does it exist ?
- Why we are interested ?
- How to observe ?
- Our CMB observation in China AliCPT

What is CMB?

- Cosmic Microwave Background radiation (CMB)
 - Relic density from Big Bang
 - CMB is an old fossil !





CMB was discovered in 1965



What a Discovery!

In view of the similarities, Penzias and Wilson began to realize what they had encountered. The noise that had been perplexing them was actually the cosmic background radiation (CMB).



Why we are interested?

Cosmic Microwave Background Spectrum from CCBE









> One hundred thousandths temperature fluctuation $\frac{\Delta T}{T} \sim 10^{-5}$

Observation of CMB



Old fossil: we trace back to t~380,000 yrs



Cosmic evolution between: t of [380,000yrs, 138Gyrs]



A Precision Cosmological Probe

• Trace back to very early: last scattering

 $t \sim 380,000 \ yrs, \qquad T \sim 1 eV$

- Tracer is simple and clean: free streaming after decoupling
- Framework is simple: thermal dynamics + linear perturbation theory

$\frac{\Delta T}{T}$: what can we learn ?

 \blacktriangleright Early universe: inflation, bouncing...

Geometry: Curvature

Energy component: baryon, DM, DE…

≻Late time evolution: cosmic expansion







• The Guidepost





or

Background evolution: expansion

• Friedmann-Roberston-Walker (FRW) spacetime

$$ds^{2} = dt^{2} - a^{2}(t)\left[\frac{dr^{2}}{1 - kr^{2}} + r^{2}d\theta^{2} + r^{2}\sin^{2}\theta d\varphi^{2}\right]$$

$$ds^2 = a^2(\eta)(d\eta^2 - \gamma_{ij}dx^i dx^j)$$

 η is conformal time

$$\eta = \int dt/a$$

 $x^1 = R \sin \theta \cos \varphi, \ x^2 = x^1 = R \sin \theta \sin \varphi, \ x^3 = R \cos \theta$

$$\gamma_{ij} = \frac{\delta_{ij}}{(1 + \frac{K}{4}x^k x^k)^2} \qquad r = \frac{R}{(1 + \frac{K}{4}R^2)^2}$$

Comoving:
$$\dot{r}=\dot{ heta}=\dot{arphi}=0$$



Scale factor a(t)

Cosmic evolution: background

• GR equation: dynamics of universe

$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} = -8\pi G T_{\mu\nu}$$

Perfect fluid $T_{\mu\nu} = (\rho + p)U_{\mu}U_{\nu} - pg_{\mu\nu}$ $U_{\mu}U^{\mu} = 1$

$$U^i = 0, \ U_i = 0$$

| with cosmic time t | with conformal time η |
|--|--|
| $U^0 = 1, \ U_0 = 1$ | $U^0 = 1/a, \ U_0 = a$ |
| $\begin{split} H^2 + \frac{K}{a^2} &= \frac{8\pi G}{3}\rho\\ \frac{\ddot{a}}{a} &= -\frac{4\pi G}{3}(\rho + 3p) \end{split}$ | $\mathcal{H}^2 + K = \frac{8\pi G}{3}a^2\rho$ $\mathcal{H}' = -\frac{4\pi G}{3}a^2(\rho + 3p)$ |
| $\dot{\rho} + 3H(\rho + p) = 0$ | $\rho' + 3\mathcal{H}(\rho + p) = 0$ |

 $\omega = p/\rho$ Equation of state is required

Multi components

$$\rho = \sum_{i} \rho_i, \ p = \sum_{i} p_i$$

$$1 + \frac{k}{a^2 H^2} = \frac{8\pi G}{3H^2} \rho \equiv \Omega \qquad \qquad \Omega = \sum \Omega_i$$

$$\Omega_k \equiv -\frac{k}{a^2 H^2} \qquad \qquad \Omega + \Omega_k = 1$$



$$\begin{split} \dot{\rho}_i + 3H(\rho_i + p_i) = 0 \\ \hline{\rho}_i + 3H(\rho_i + p_i) = 0 \\ \hline{P}_i = p_i(\rho_i) \\ With constant EoS \\ Radiation \\ Radiation \\ w_r = 1/3, \rho_r \propto a^{-4} \\ Matter (non-relativistic) \\ W_n = -1, \rho_A = constant \\ \end{split}$$

Single component universe (K=0)

 $a \propto t^{rac{2}{3(1+w)}}$ or $a \propto \eta^{rac{2}{1+3w}}$ Radiation domination $a \propto t^{rac{1}{2}}$ $a \propto \eta$

Matter domination $a \propto t^{rac{2}{3}}$ $a \propto \eta^2$



Linear perturbaton: Cosmic fluctuation



$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

$$\delta G_{\mu\nu} = 8\pi G \delta T_{\mu\nu}$$

(spacetime)

Spacetime perturbations

 $ds^{2} = a(\eta)^{2} \{ (1+2A)d\eta^{2} - 2(B_{|i} + S_{i})d\eta dx^{i} - [(1-2\psi)\gamma_{ij} + 2E_{|ij} + 2F_{(i|j)} + h_{ij}] dx^{i} dx^{j} \}$ $S_{|i}^{i} = F_{|i}^{i} = 0, \ h_{i}^{i} = 0, \ h_{j|i}^{i} = 0$

Gauge transformations and gauge invariant perturbations

Scalar

$$A \to \tilde{A} = A - \mathcal{H}\xi^{0} - \xi^{0'}$$

$$\psi \to \tilde{\psi} = \psi + \mathcal{H}\xi^{0}$$

$$B \to \tilde{B} = B + \xi^{0} - \xi'$$

$$E \to \tilde{E} = E - \xi$$

$$\Phi = A + (1/a)[(B - E')a]',$$

$$\Psi = \psi - \mathcal{H}(B - E');$$

Vector

Tensor $\tilde{h}_{ij} = h_{ij}$

Matter perturbations

 $T^{\mu}_{\nu}(\eta,\vec{x}) = T^{\mu}_{\nu}(\eta) + \delta T^{\mu}_{\nu}(\eta,\vec{x})$

 $\delta \tilde{T}^{\mu}_{\nu} = \delta T^{\mu}_{\nu} - \mathcal{L}_{\xi} T^{\mu}_{\nu}$

(matter)

Matter as a fluid

 $T^{\mu}_{\nu} = -p\delta^{\mu}_{\nu} + (\rho + p)U^{\mu}U_{\nu} + \Sigma^{\mu}_{\nu} \qquad U^{\mu}U_{\mu} = 1 \qquad \Sigma_{\mu\nu}U^{\nu} = 0 , \quad \Sigma^{\mu}_{\mu} = 0$ On background $\rho, p, U^{0} = 1/a, U^{i} = 0$

> anisotropic stress due to shear viscosity and heat flow, vanishes on background, gauge invariant perturbations

At linear order

$$\begin{split} \delta\rho, \delta p, \quad \delta U^0 &= -\frac{A}{a} , \ \delta U_0 = aA \\ U_i &= U_{|i} + U_i^{vec} , \ U_i^{vec|i} = 0 \\ \Sigma_{00} &= 0 , \ \Sigma_{0i} = 0 , \ \Sigma_i^i = 0 \\ \Sigma_{ij} &= -a^2 [\Sigma_{|ij} - \frac{1}{2} \nabla^2 \Sigma \gamma_{ij} + \Sigma_{(i|j)} + \sigma_{ij}] \end{split}$$

Scalar

$$\begin{split} \tilde{\delta\rho} &= \delta\rho - \rho'\xi^0 & \tilde{U} = U - a\xi^0 \\ \tilde{\delta p} &= \delta p - p'\xi^0 & \tilde{\Sigma} = \Sigma \end{split}$$

Vector
$$\tilde{U}_i^{vec} = U_i^{vec}$$
 $\tilde{\Sigma}_i = \Sigma_i$

Tensor
$$\tilde{\sigma}_{ij} = \sigma_{ij}$$

Some gauge invariant perturbations

$$^{(gi)}\delta\rho = \delta\rho + \rho'(B - E') , \ ^{(gi)}\deltap = \delta p + p'(B - E') , \ ^{(gi)}U = U + a(B - E')$$
$$\mathcal{R} = -\psi - \mathcal{H}\frac{\delta\rho}{\rho'} , \ \delta p_{nad} \equiv \delta p - \frac{p'}{\rho'}\delta\rho$$

Vector perturbation decay gradually

Vector perturbation

$$k^2 \mathcal{F}_i = 16\pi Ga(\rho + p)U_i^{vec} ,$$

$$(\mathcal{F}_{i,j} + \mathcal{F}_{j,i})' + 2\mathcal{H}(\mathcal{F}_{i,j} + \mathcal{F}_{j,i}) = 0$$

$$\mathcal{F}_i \propto 1/a^2 \qquad \quad U_i^{vec} \propto 1/[a^3(\rho+p)]$$

$$V_i = V^i = \frac{U_i^{vec}}{a^4(\rho + p)}$$

在辐射为主时期 ρ , $p \propto a^{-4}$, 因此 $V_i = const.$; 在物质为主时期 ρ , $p \propto a^{-3}$, 因此 $V_i \propto \frac{1}{a}$ 是逐渐衰减

Tensor perturbation: Gravitational wave

$$h_{ij}'' + 2\mathcal{H}h_{ij}' + k^2h_{ij} = 16\pi Ga^2\sigma_{ij}$$

Scalar perturbation

$$\begin{split} ds^2 &= a^2 [(1+2\Phi) d\eta^2 - (1-2\Psi) \delta_{ij} dx^i dx^j] \\ & 3\mathcal{H}(\mathcal{H}\Phi + \Psi') + k^2 \Psi = -4\pi G a^2 \delta \rho \ , \\ & \mathcal{H}\Phi + \Psi' = 4\pi G a (\rho + p) U \ , \\ & (2\mathcal{H}' + \mathcal{H}^2) \Phi + \mathcal{H}\Phi' + \Psi'' + 2\mathcal{H}\Psi' + \frac{k^2}{3} (\Psi - \Phi) = 4\pi G a^2 \delta p \\ & k^2 (\Psi - \Phi) = 12\pi G a^2 (\rho + p) \sigma \ . \end{split}$$

 $\Phi'' + 3\mathcal{H}(1 + c_s^2)\Phi' - c_s^2\nabla^2\Phi + [2\mathcal{H}' + (1 + 3c_s^2)(\mathcal{H}^2 - K)\Phi] = 4\pi Ga^2\delta p_{nad}$



Coupled system: scattering & interaction



$$C[f] = \int d(\text{phase space})[\text{ energy-momentum conservation}]$$

 $\times |M|^2[\text{emission} - \text{absorption}]$

 $f(x^i, P_i, t)$ distribution function $dx^1 dx^2 dx^3 dP_1 dP_2 dP_3$ $dN = f dx^1 dx^2 dx^3 dP_1 dP_2 dP_3$ 能动量张量 $T^{\mu}_{\nu} = \frac{g}{(2\pi\hbar)^3} \int \frac{dP_1 dP_2 dP_3}{\sqrt{-\det|q|}P^0} P^{\mu} P_{\nu} f(x^i, P_j, t)$ $\rho = T_0^0 = \frac{g}{(2\pi)^3} \int d^3 p E f(\vec{x}, \vec{p}, t)$ In FRW: $T_{j}^{i} = -P\delta_{j}^{i} = \frac{g}{(2\pi)^{3}} \int \frac{d^{3}p}{E} (-p^{i}p^{j})f(\vec{x}, \vec{p}, t)$ $= -\frac{g}{(2\pi)^3} \int \frac{d^3p}{E} \frac{p^2}{3} f(\vec{x}, p, t) \delta_j^i$ $P = \frac{g}{(2\pi)^3} \int \frac{d^3p}{E} \frac{p^2}{3} f(\vec{x}, p, t)$ $n = \frac{g}{(2\pi)^3} \int d^3 p f(\vec{x}, \vec{p}, t)$

In thermal equilibrium:

$$f(\vec{x}, \vec{p}, t) = \frac{1}{\exp \frac{E - \mu}{T} \pm 1}$$

Perturbed BE

Newtonian gauge: $ds^2 = a^2(1+2\phi)d\eta^2 - a^2[(1-2\psi)\delta_{ij} + h_{ij}]dx^i dx^j$ $f(x^i, P_j, \eta) = f(x^i, q, n_j, \eta) = f_0(q, \eta) + \delta f(x^i, q, n_j, \eta)$ Photon: $f(x^i, n_j, \eta) = [exp\{\frac{p}{T(\eta)[1 + \Delta(x^i, n_j, \eta)]}\} - 1]^{-1}$

BE:
$$\frac{\partial f}{\partial \eta} + \frac{dx^i}{d\eta} \frac{\partial f}{\partial x^i} + \frac{dq}{d\eta} \frac{\partial f}{\partial q} + \frac{dn_i}{d\eta} \frac{\partial f}{\partial n_i} = (\frac{\partial f}{\partial \eta})_C$$

$$\frac{dx^i}{d\eta} = \frac{P^i}{P^0} = -\frac{q}{\epsilon}(1+\phi+\psi)n_i$$

Geodesic equation

$$P^{0}\frac{dP^{\mu}}{d\eta} + \Gamma^{\mu}_{\alpha\beta}P^{\alpha}P^{\beta} = 0 \quad \Longrightarrow \quad \frac{dq}{d\eta} = q\psi' - \epsilon n_{i}\partial_{i}\phi$$

Perturbed Boltzmann equation

$$\frac{\partial \delta f}{\partial \eta} + i \frac{q}{\epsilon} k \mu \ \delta f + q \frac{\partial f_0}{\partial q} (\psi' - i \frac{\epsilon}{q} k \mu \phi) = (\frac{\partial f}{\partial \eta})_C$$

Equations of the system: 9

$$\frac{\partial \Delta_T^{(S)}}{\partial \eta} + ik\mu \Delta_T^{(S)} = \psi' - ik\mu\phi + k[-\Delta_T^{(S)} + \Delta_{T0}^{(S)} + i\mu v_b + \frac{1}{2}P_2(\mu)\Pi]$$
(1)

• photons: $\frac{\partial \Delta_P^{(S)}}{\partial \eta} + ik\mu \Delta_P^{(S)} = k\{-\Delta_P^{(S)} + \frac{1}{2}[1 - P_2(\mu)]\Pi\}$ (2)

$$\Pi = \Delta_{T2}^{(S)} + \Delta_{P2}^{(S)} + \Delta_{P0}^{(S)} \quad \dot{\kappa} = an_e X_e \sigma_T \quad 1 - P_2(\mu) = \frac{3}{2}(1 - \mu^2)$$
(3)

$$\delta_b' = -kv_b + 3\psi' \tag{5}$$

(8)

$$v'_{b} = -\mathcal{H}v_{b} + c_{s}^{2}k\delta_{b} + \frac{4\rho_{\gamma}}{3\rho_{b}}\dot{\kappa}(3\Delta_{T1}^{(S)} - v_{b}) + k\phi \qquad c_{s}^{2} = \frac{P'_{b}}{\rho'_{b}} \quad (6)$$

• Neutrino:
$$\frac{\partial F_{\nu}}{\partial \eta} + ik\mu F_{\nu} = 4(\psi' - ik\mu\phi)$$
 (7)

• Einstein EQ:
$$F_{\nu}(\vec{h}, \hat{n}, \eta) = \frac{\int dq q^3 \, \delta f}{\int dq q^3 f_0} = \sum_l (-i)^l (2l+1) F_{\nu l}(\vec{h}, \eta) P_l(\mu)$$

• Baryon:

evolution

$$k(\psi' + \mathcal{H}\phi) = 4\pi Ga^2 \sum (\rho_i + P_i)v_i$$
(9)

Linear
Boltzmann
Initial
inhomogeneity
Perturbation
Fas
condition

Initial condition: Inflation

- Horizon problem
- Flatness problem





Initial condition



• Initial condition from inflation:

 $\left\langle \Phi(\vec{k}) \right\rangle = 0$

Linear perturbation

 $\langle \Phi(\vec{k}) \Phi^*(\vec{k'}) \rangle = (2\pi)^3 P_{\Phi}(k) \delta^3(\vec{k} - \vec{k'})$



temperature



polarization

Initial condition

• Inflation tell you : primordial power spectrum

$$k^{3} P_{\Phi}(k) = \frac{8\pi}{9} \frac{H^{2}}{\epsilon m_{pl}^{2}} \Big|_{aH=k} \equiv A_{S}(\frac{k}{k_{*}})^{n_{S}-1}$$

$$r \equiv k^{3} P_{h}(k) = 8\pi \frac{H^{2}}{m_{pl}^{2}} \Big|_{aH=k} \equiv A_{T}(\frac{k}{k_{*}})^{n_{T}}$$

• Adiabatic conditions: $\delta p = c_s^2 \delta \rho$, ignore : δS



Equations of the system: 9

$$\frac{\partial \Delta_T^{(S)}}{\partial \eta} + ik\mu \Delta_T^{(S)} = \psi' - ik\mu\phi + \kappa \left[-\Delta_T^{(S)} + \Delta_{T0}^{(S)} + i\mu v_b + \frac{1}{2}P_2(\mu)\Pi \right]$$
(1)

• photons:

$$\frac{\partial \Delta_P^{(S)}}{\partial \eta} + ik\mu \Delta_P^{(S)} = k\{-\Delta_P^{(S)} + \frac{1}{2}[1 - P_2(\mu)]\Pi\} \qquad (2)$$

$$\Pi = \Delta_{T2}^{(S)} + \Delta_{P2}^{(S)} + \Delta_{P0}^{(S)} \quad \dot{\kappa} = an_e X_e \sigma_T \quad 1 - P_2(\mu) = \frac{3}{2}(1 - \mu^2)$$

$$\delta_b' = -kv_b + 3\psi' \tag{5}$$

Baryon:

$$v'_{b} = -\mathcal{H}v_{b} + c_{s}^{2}k\delta_{b} + \frac{4\rho_{\gamma}}{3\rho_{b}}\dot{\kappa}(3\Delta_{T1}^{(S)} - v_{b}) + k\phi$$
 $c_{s}^{2} = \frac{P'_{b}}{\rho'_{b}}$ (6)

• Neutrino:
$$\frac{\partial F_{\nu}}{\partial \eta} + ik\mu F_{\nu} = 4(\psi' - ik\mu\phi)$$
 (7)

background

evolution

Perturbation

• Einstein EQ:
$$F_{\nu}(\vec{k}, \hat{n}, \eta) = \frac{\int dq q^3 \,\delta f}{\int dq q^3 f_0} = \sum_l (-i)^l (2l+1) F_{\nu l}(\vec{k}, \eta) P_l(\mu)$$
(8)

$$k(\psi' + \mathcal{H}\phi) = 4\pi G a^2 \sum (\rho_i + P_i) v_i \tag{9}$$

condition

inhomogeneity

Focus on : $\frac{\Delta T}{T}$

• Solve BE: multipole expansion, ignore higher order, just keep order of 0 & 1

$$\begin{split} \frac{\partial \Delta_{T0}^{(S)}}{\partial \eta} &= -k \Delta_{T1}^{(S)} + \psi' \\ \frac{\partial \Delta_{T1}^{(S)}}{\partial \eta} &= \frac{k}{3} (\Delta_{T0}^{(S)} - 2\Delta_{T2}^{(S)} + \phi) + \dot{\kappa} (\frac{v_b}{3} - \Delta_{T1}^{(S)}) \\ \frac{\partial \Delta_{T2}^{(S)}}{\partial \eta} &= \frac{k}{5} (2\Delta_{T1}^{(S)} - 3\Delta_{T3}^{(S)}) + \dot{\kappa} (\frac{\Pi}{10} - \Delta_{T2}^{(S)}) \\ \frac{\partial \Delta_{Tl}^{(S)}}{\partial \eta} &= \frac{k}{2l+1} [l \Delta_{T(l-1)}^{(S)} - (l+1) \Delta_{T(l+1)}^{(S)}] - \dot{\kappa} \Delta_{Tl}^{(S)}, \ l \ge 3 \\ \frac{\partial \Delta_{Pl}^{(S)}}{\partial \eta} &= \frac{k}{2l+1} [l \Delta_{P(l-1)}^{(S)} - (l+1) \Delta_{P(l+1)}^{(S)}] + \dot{\kappa} [-\Delta_{Pl}^{(S)} + \frac{1}{2} \Pi (\delta_{l0} + \frac{\delta_{l2}}{5})] \end{split}$$

For strong coupling $(\tau \gg 1)$: $\Delta_l \sim \frac{k\eta}{2\tau} \Delta_{l-1}$

1995: Wayne Hu& ...

 $\Delta_0(\eta) + \psi(\eta) \sim [\Delta_0(0) + \psi(0)] \cos(kr_s)$

 $\Delta_1(\eta) \sim [\Delta_0(0) + \psi(0)] \sin(kr_s)$

Oscillate



From Δ_T to angular power spectrum:

- what inflation predict is the variance of Φ

 $\langle \Phi(\vec{k}) \Phi^*(\vec{k'}) \rangle = (2\pi)^3 P_{\Phi}(k) \delta^3(\vec{k} - \vec{k'})$

• It's useful to do the decomposition into harmonics

$$\Delta_T(\hat{n}) = \sum_{lm} a_{lm} Y_{lm}(\hat{n})$$

• angular power spectrum

$$< a_{lm}a_{lm}^* > = C_l \delta_{ll'} \delta_{mm'}$$

$$C_l = \delta_{ll'} \delta_{mm'} 4\pi \int dlnk j_l^2 (kD_*) \frac{k^3}{2\pi^2} P_{\Phi/h}(k)$$



TT power spectrum:









声学振荡:

早期宇宙温度很高,处于等离子体状态,粒子之间相互耦合。扰动以声学振荡的 形式在此等离子体中传播,直到光子与重子物质退耦合形成CMB。

声学视界:即到再复合(Recombination)时,声波能够传播的距离



from Wayne Hu



First peak position:

curvature of universe



$$\sum \Omega_i + \Omega_k = 1$$



from Wayne Hu

CMB中的宇宙学信息



from Wayne Hu
CMB中的宇宙学信息

The difference between 2nd and 3rd peaks: baryon !

• Baryon add extra mass to the rbaryon fluid, the controlling parameter is the momentum density ratio

$$R \equiv \frac{p_b + \rho_b}{p_\gamma + \rho_\gamma} \approx \frac{30\Omega_b h^2}{10^{-3}} \left(\frac{a}{10^{-3}}\right)$$

Which is important at recombination ! Of order unity

Modification to the solution

 $[\Delta_T + (1 + \mathbf{R})\Psi](\eta) = [\Delta_T + (1 + \mathbf{R})\Psi](0)\cos(\theta)$

Even-odd peak modulation of effective temperature

$$egin{aligned} [\Theta+\Psi]_{ ext{peaks}} &= [\pm(1+3R)-3R]\,rac{1}{3}\Psi(0)\ \Theta+\Psi]_1 - [\Theta+\Psi]_2 &= [-6R]rac{1}{3}\Psi(0) \end{aligned}$$



CMB中的宇宙学信息

Damping tail : baryon



from Wayne Hu

The latest measurements



On inflation



The constraints

1807.06209

Table 1

| Parameter | TT+lowE 68% limits | TE+lowE 68% limits | EE+lowE 68% limits | TT,TE,EE+lowE 68% limits | TT,TE,EE+lowE+lensing 68% limits | TT,TE,EE+lowE+lensing+BAO 68% limits |
|---|-----------------------|-----------------------|------------------------|------------------------------|-------------------------------------|---|
| Ω ₂ μ ² | 0.02212 ± 0.00022 | 0.02249 ± 0.00025 | 0.0240 ± 0.0012 | 0.02236 ± 0.00015 | 0.02237 ± 0.00015 | 0.02242 ± 0.00014 |
| $\Omega_c h^2$ | 0.1206 ± 0.0021 | 0.1177 ± 0.0020 | 0.1158 ± 0.0046 | 0.1202 ± 0.0014 | 0.1200 ± 0.0012 | 0.11933 ± 0.00091 |
| 1039 _{NC} | 1.04077 ± 0.00047 | 1.04139 ± 0.00049 | 1.03999 ± 0.00089 | 1.04090 ± 0.00031 | 1.04092 ± 0.00031 | 1.04101 ± 0.00029 |
| T + + + + + + + + + + + + + + + + + + + | 0.0522 ± 0.0080 | 0.0496 = 0.0085 | 0.0527 ± 0.0090 | $0.0544^{+0.0010}_{-0.0051}$ | 0.0544 = 0.0073 | 0.0561 ± 0.0071 |
| $ln(10^{10}A_s)\ldots\ldots,$ | 3.040 ± 0.016 | 3.018+0.000 | 3.052 ± 0.022 | 3.045 ± 0.016 | 3.044 ± 0.014 | 3.047 ± 0.014 |
| <i>R</i> ₂ | 0.9626 ± 0.0057 | 0.967 ± 0.011 | 0.980 ± 0.015 | 0.9649 ± 0.0044 | 0.9649 ± 0.0042 | 0.9665 ± 0.0038 |
| $H_0 [{\rm kms^{-1}Mpc^{-1}}]$ | 66.88 ± 0.92 | 68.44 ± 0.91 | 699127 | 67.27 ± 0.60 | 67.36 ± 0.54 | 67.66 ± 0.42 |
| Ω | 0.679 ± 0.013 | 0.699 ± 0.012 | 0.711+0.025 | 0.6834 ± 0.0084 | 0.6847 = 0.0073 | 0.6889 ± 0.0056 |
| Ω _m | 0.321 ± 0.013 | 0.301 ± 0.012 | 0.289+0/05 | 0.3166 ± 0.0084 | 0.3153 = 0.0073 | 0.3111 ± 0.0056 |
| $\Omega_{\pi}h^2$ | 0.1434 ± 0.0020 | 0.1408 ± 0.0019 | 0.1404_00054 | 0.1432 ± 0.0013 | 0.1430 ± 0.0011 | 0.14240 ± 0.00087 |
| Ω ₃₁ ^{β³} | 0.09589 + 0.00046 | 0.09635 ± 0.00051 | 0.0981_00018 | 0.09633 ± 0.00029 | 0.09633 ± 0.00030 | 0.09635 ± 0.00030 |
| σι | 0.8118 ± 0.0089 | 0.793 ± 0.011 | 0.796 ± 0.018 | 0.8120 ± 0.0073 | 0.8111 = 0.0060 | 0.8102 ± 0.0060 |
| $S_{0}=\sigma_{0}(\Omega_{m}/0.3)^{0.5} .$ | 0.840 ± 0.024 | 0.794 ± 0.024 | 0.781-00% | 0.834 ± 0.016 | 0.832 ± 0.013 | 0.825 ± 0.011 |
| $\sigma_{\rm F} \Omega_{\rm m}^{129}$ | 0.611 ± 0.012 | 0.587 ± 0.012 | 0.583 ± 0.027 | 0.6090 ± 0.0081 | 0.6078 ± 0.0064 | 0.6051 ± 0.0058 |
| In | 7.50 ± 0.82 | 7.11 (0.8) | 7.10 0.37 | 7.68 ± 0.79 | 7.67 ± 0.73 | 7.82 ± 0.71 |
| 10 ⁹ A, | 2.092 ± 0.034 | 2.045 ± 0.041 | 2.116 ± 0.047 | $2.101_{-0.044}^{+0.081}$ | $2,100 \pm 0.030$ | 2.105 ± 0.030 |
| 10 ^a A ₂ e ⁻²⁷ | 1.884 ± 0.014 | 1.851 ± 0.018 | 1.904 ± 0.024 | 1.884 = 0.012 | 1.883 ± 0.011 | 1.881 ± 0.010 |
| Age Gyr | 13.830 ± 0.037 | 13.761 ± 0.038 | 13.64+0.16 | 13.800 ± 0.024 | 13.797 = 0.023 | 13.787 ± 0.020 |
| z | 1090.30 ± 0.41 | 1089.57 ± 0.42 | $1087.8^{-1.6}_{-1.2}$ | 1089.95 ± 0.27 | 1089.92 = 0.25 | 1089.80 ± 0.21 |
| r. [Mpc] | 144.46 ± 0.48 | 144.95 ± 0.48 | 144.29 ± 0.64 | 144.39 ± 0.30 | 144.43 ± 0.26 | 144.57 ± 0.22 |
| 1009 | 1.04097 ± 0.00046 | 1.04156 ± 0.00049 | 1.04001 ± 0.00086 | 1.04109 ± 0.00030 | 1.04110 ± 0.00031 | 1.04119 ± 0.00029 |
| 2 ₁₁₄ | 1059.39 ± 0.46 | 1060.03 ± 0.54 | 1063.2 ± 2.4 | 1059.93 ± 0.30 | 1059.94 ± 0.30 | 1060.01 ± 0.29 |
| r _{d vg} [Mpc] | 147.21 ± 0.48 | 147.59 ± 0.49 | 146.46 ± 0.70 | 147.05 ± 0.30 | 147.09 = 0.26 | 147.21 ± 0.23 |
| $h_{\rm D}$ [Mpc ⁻¹] | 0.14054 ± 0.00052 | 0.14043 ± 0.00057 | 0.1426 ± 0.0012 | 0.14090 ± 0.00032 | 0.14087 ± 0.00030 | 0.14078 ± 0.00028 |
| 2.4. · · · · · · · · · · · · | 3411 + 48 | 3349 ± 46 | 3340_61 | 3407 ± 31 | 3402 ± 26 | 3387 ± 21 |
| λ _{eq} [Mpc ⁻¹] | 0.01041 ± 0.00014 | 0.01022 ± 0.00014 | 0.01019-0.903 | 0.010398 ± 0.000094 | 0.010384 ± 0.000081 | 0.010339 ± 0.000063 |
| 1009 _{8,54} | 0.4483 ± 0.0046 | 0.4547 ± 0.0045 | 0.4562 ± 0.0092 | 0.4490 ± 0.0030 | 0.4494 ± 0.0026 | 0.4509 ± 0.0020 |

CMB polarization:

 \succ Thomson scattering can only form linear polarization, so V = 0.





Stokes parameters

Description

$$E_x = a_x(t) \cos[\omega_0 t - \theta_x(t)] \quad I = \langle a_x^2 \rangle + \langle a_y^2 \rangle, \quad U = \langle 2a_x a_y \cos(\theta_x - \theta_y) \rangle,$$

 $E_{v} = a_{y}(t) \cos[\omega_{0}t - \theta_{y}(t)] \quad Q = \langle a_{x}^{2} \rangle - \langle a_{y}^{2} \rangle, \quad V = \langle 2a_{x}a_{y}\sin(\theta_{x} - \theta_{y}) \rangle.$

I: intensity , *Q* and *U: linear polarization*, *V*: circular polarization

Elliptical polarization is the combination of linear polarization and circular polarization, which is the most general case of photon polarization.





Coordinate independent quantities

Q, U are not coordinate invariants

$$Q \pm iU \to e^{\mp 2i\alpha} (Q \pm iU)$$

E,B decomposition:

$$(Q+iU)(\hat{n}) = \sum_{lm} a_{2,lm2} Y_{lm}(\hat{n})$$
$$(Q-iU)(\hat{n}) = \sum_{lm} a_{-2,lm-2} Y_{lm}(\hat{n})$$
$$a_{E,lm} = -(a_{2,lm} + a_{-2,lm})/2,$$

 $a_{B,lm} = i(a_{2,lm} - a_{-2,lm})/2.$

$$\begin{split} a_{2,lm} &= \int d\Omega_2 Y^*_{lm}(\hat{\pmb{n}}) (\underline{Q} + iU)(\hat{\pmb{n}}) \\ &= \left[\frac{(l+2)!}{(l-2)!} \right]^{-1/2} \int d\Omega Y^*_{lm}(\hat{\pmb{n}}) \delta^2 (\underline{Q} - iU)(\hat{\pmb{n}}), \\ a_{-2,lm} &= \int d\Omega_{-2} Y^*_{lm}(\hat{\pmb{n}}) (\underline{Q} - iU)(\hat{\pmb{n}}) \\ &= \left[\frac{(l+2)!}{(l-2)!} \right]^{-1/2} \int d\Omega Y^*_{lm}(\hat{\pmb{n}}) \delta^2 (\underline{Q} - iU)(\hat{\pmb{n}}). \end{split}$$

E: parity even, B: parity odd E-mode B mode

Map of E & B



E-mode

B-mode

R. Durrer

$$\begin{split} T(\hat{n}) &= \sum_{lm} a_{T,lm} Y_{lm}(\hat{n}) \\ (Q+iU)(\hat{n}) &= \sum_{lm} a_{2,lm2} Y_{lm}(\hat{n}) \\ (Q-iU)(\hat{n}) &= \sum_{lm} a_{-2,lm-2} Y_{lm}(\hat{n}) \\ < a_{T,l'm'}^{*} a_{T,lm} > = C_{l}^{TT} \delta_{ll'} \delta_{mm'} \\ < B_{l'm'}^{*} B_{lm} > = C_{l}^{TB} \delta_{ll'} \delta_{mm'} \\ < a_{T,l'm'}^{*} B_{lm} > = C_{l}^{TB} \delta_{ll'} \delta_{mm'} \\ < a_{T,l'm'}^{*} B_{lm} > = C_{l}^{TB} \delta_{ll'} \delta_{mm'} \\ < B_{l'm'}^{*} B_{lm} > = C_{l}^{TB} \delta_{ll'} \delta_{mm'} \\ < B_{l'm'}^{*} B_{lm} > = C_{l}^{TB} \delta_{ll'} \delta_{mm'} \\ < B_{l'm'}^{*} B_{lm} > = C_{l}^{TB} \delta_{ll'} \delta_{mm'} \\ \end{cases}$$

B mode: primordial gravitational wave

Scalar quadrupole, azimuthal symmetric generates T and E



Tensor quadrupole, without azimuthal symmetry, generates T, E and B





Tensile & compression







The next goal: Primordial gravitational waves has not been discovered yet !



The detection

CMB photon number density

400–500 photons/cm³

How weak it is?



- Peak wavelength : ג=2mm
- CMB Radiation Intensity : 3.3 microwatts/m²
- Radiation intensity of human body : 500 watts/m²

How to detect ?



Photons get absorbed by block, Block heats up, Thermometer senses temperature rise.

How to observe ?

• Want



1. Super-sensitive thermometer

 Small, very cold, block (so CMB photons heat it up fast)

3. No "extra" photons from our hot environment.

Sensitive Thermometers: TES

AlMn

film

- Optical load :
 - $P_{load} = 2\eta k T_{RJ} \Delta v + P_{internal}$
- Saturation power:
 - $P_{saturation} = T_0 G_0 \frac{(T_c/T_0)^{n+1}-1}{n+1} \sim 2P_{load}$
- Detector noise: $NEP = \sqrt{NEP_{photon}^2 + NEP_{phonon}^2 + NEP_{other}^2}$
 - Photon noise: $NEP_{photon} = \sqrt{2hvP_{load} + \frac{2P_{load}^2}{v(\Delta v/v)}}$
 - Phonon noise : $NEP_{phonon} = \sqrt{4k_BGT^2F}, \quad G = G_0 \left(\frac{T}{T_0}\right)^n$
 - others







CMB photon detector module



Arxiv: 1607.06861, BICEP3 focal plane design and detector performance, BICEP collaboration

SPT Polarization "Camera"



"Pixels" sensitive to 2mm light at the center, and 3mm light (outer set)

Current "state of the art" is around 1000 pixel cameras.

Here you see the "feed horns" that funnel light to the detectors...

From SPT collaboration

SPTpol 150GHz Detectors (built in a cleanroom at NIST)



(A few years to develop the ability to make a working first copy, ~1 month for a batch of 2-3 after that. SPTpol has 7 copies.)

From SPT collaboration

Polarization measurements: experiments







Measurements

- WMAP satellite has mapped the temperature fluctuations over the entire sky, the region around the third peak has been filled in by balloon and groundbased experiments such as BOOMERANG and CBI, and on small scales, ACBAR recently resolved the fourth and fifth acoustic peaks for the first time.
- BOOMERANG: flat universe



Boomerang collaboration. PRL86 (2001) 3475-3479



J. Kovac, et al. nature420 (2002), 772

A new stage for BB measurement after BICEP2



- The most accurate measurement of polarization with error bars on BB Spectrum
- Component separation for foreground reduction: dust & syncrotron

What does Planck tell us ?



- multi-frequency measurements: 9 bands from 30 857 GHz
- precision measurement of T & E
- Contamination from dust can not be ignored !

The polarization detection after Planck

- Ground observations will play important role :
 - Litebird may happen in almost 10 years.
 - Ground observation have achieved large stride !



- Ground telescope have many advantages
 - a large amount of TES on the focal plane
 - Array of Telescope is easy

Ground and space are complementary





CMB S4:





CMB S4



CMB S4 white paper

CMB physics after Planck

 Searching for primordial gravitational waves: constrain the energy scale of inflation and to test alternative models, and to provide insights into quantum gravity;

Non-r

Measuring effective number of light relativistic species (dark radiation): to search for new light relics, independent tests of BBN and understand the evolution of the Universe at t = 1 sec;

- Sum of the neutrino masses;
- Dark energy study: using secondary CMB anisotropy through its impact on the growth of structure;
- Testing general relativity and constraining alternate theories of gravity on large scales.

CMB S4 white paper

Large + small aperature telescopes:

- More detectors
- Large sky coverage
- More deep obs





- need sensitivity :
 1 uk arcmin
- over half sky,

 four years survey with 500,000
 CMB-sensitive detectors;

- More sensitive detectors
- Extend to
 northern sky,
- New telescope: Current telescopes are saturated;

地面CMB探测对台址要求苛刻

• 微波波段大气透射率



地球上有四个台址



H. Li et al arXiv:1710.03047

主要集中于南半球,阿里是北半球唯一




PWV: Datasets

Y.Li, Y.Liu, S.Li, H.Li and X.Zhang, ArXiv:1709.09053

- Datasets :
 - MERRA-2 Reanalysis data (NASA) :
 - Spatial Resolution: 0.6251on*0.51at*721ayers
 - Time Resolution: 3hours
 - Data : Relative humidity, Temperature, Pressure, Altitude....
 - Radiosonde data from Ali local weather station :
 - Send the balloon twice a day : 07:00 & 19:00
 - Data : Dew-point temperature, Temperature, Pressure....

Formula:

$$PWV = \int \rho_{\nu} dh = \int q_{\nu} \rho dh = -1/g \int q_{\nu} dp \approx -1/g \sum_{i} q_{\nu}^{i} \triangle p_{i}$$



Radiosonde

Q.Ye, M.Su, H.Li, X.Zhang,ArXiv:1512.01099Chao-lin Kuo,ArXiv:1707.08400

Results: (1mm/5250m, 0.6mm/6000m)



PWV distribution (MERRA-2)

- Strong seasonal variation.
- Oct. to Mar: The median value of PWV is Radiosonde: 0.92mm/0.56mm MERRA-2: 1.07mm/0.62mm
- Results derived from satellite and radiosonde are consistent.





Comparison between MERRA-2 & radiosonde (Oct.-Mar) • Observing season: Oct. to Mar.



CMB observation in Tibet



项目概况:由中科院高能所牵头的国际合 作项目。在西藏阿里地区海拔5250米的B1 点建设高灵敏的原初引力波探测望远镜 (阿里一号,AliCPT-1)。





2. 时间表:

2014年5月提出; 2016年12月正式启动; 2017年3月奠基; 2018年11月观测仓验收; 预期2020-2021观测季开始观测

3. 科学目标:

瞄准原初引力波,首次给出对北半球CMB极化的 最精确测量,探寻宇宙起源。



Figure Credit: AliCPT-1 team

建成后,我国的原初引力波研究进入国际前沿 参与国际合作南北协调观测,实现国际上最灵敏的探测 阿里与南极、智利一起成为国际上CMB探测的三大基地



Scientific goal



- Probing the primordial gravitational waves (PGWs) with BB spectra.
- Measuring the rotation angle, testing CPT symmetry with TB and EB spectra.
- Investigating the CMB polarization hemispherical asymmetry.
- Studying the cross-correlation between AliCPT and DESI.
- Studying the galactic foreground.

Current status

- Finish site construction
- Test mount in factory
- Develop telescope
- Construction for control system, operation system and science analysis platform







Preliminary concern on Scan strategy

- Observation in Ali
 - Sky move fast: helpful for large sky coverage, but not easy to focus within one small patch.
 - "clean" area in north, also in south: helpful for going deep in north, and cross check in Sourth.







simulation results, Preliminary

Simulation: sensitivity

AliCPT-1 is designed to have its first light in the winter of 2020



PLANCK 100 GHz & 143 GHz Pol ~50 uK arcmin in total AliCPT 95 GHz & 150 GHz & Total 4 modules 1 season, median map-depth 14 uK arcmin

- Probing the primordial gravitational waves (PGWs) with **BB** spectra.
- Measuring the rotation angle, testing CPT symmetry with TB and EB spectra.
- Studying the cross-correlation between AliCPT and DESI.
- Studying the galactic foreground.
- Investigating the CMB polarization hemispherical asymmetry.

Scientific goal from Xinmin Zhang

• .

Preliminary results from simulation



Simulation: noise map

Simulation: BB power spectrum

Simulation: r sensitivity

Cosmological study with AliCPT-1

• Cross correlate with DESI

• Scan galaxy: signals for foreground







AliCPT-1 map depth

patch1_mapdepth,unit=uk arcmin.Ra=272.4, Dec=55.6









Observational time: 795 h Minimum map depth: 6.9 µk arcmin Maximum map depth: 13.8 µk arcmin

Planck map depth: 48 $\mu k \ arcmin$ (all channels)



